

# Green/Yellow Solid-State Lighting via Radiative and Nonradiative Energy Transfer Involving Colloidal Semiconductor Nanocrystals

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**Abstract**—LEDs made of  $\text{In}_x\text{Ga}_{1-x}\text{N}$  and  $(\text{Al}_x\text{Ga}_{1-x})_{1-y}\text{In}_y\text{P}$  suffer from significantly reduced quantum efficiency and luminous efficiency in the green/yellow spectral ranges. To address these problems, we present the design, growth, fabrication, hybridization, and characterization of proof-of-concept green/yellow hybrid LEDs that utilize radiative and nonradiative [Förster resonance energy transfer (FRET)] energy transfers in their colloidal semiconductor nanocrystals (NCs) integrated on near-UV LEDs. In our first NC-LED, we realize a color-converted LED that incorporate green-emitting CdSe/ZnS core/shell NCs ( $\lambda_{\text{PL}} = 548$  nm) on near-UV InGaN/GaN LEDs ( $\lambda_{\text{EL}} = 379$  nm). In our second NC-LED, we implement a color-converted FRET-enhanced LED. For that, we hybridize a custom-design assembly of cyan- and green-emitting CdSe/ZnS core/shell NCs ( $\lambda_{\text{PL}} = 490$  and  $548$  nm) on near-UV LEDs. Using a proper mixture of differently sized NCs, we obtain a quantum efficiency enhancement of 9% by recycling trapped excitons via FRET. With FRET-NC-LEDs, we show that it is possible to obtain a luminous efficacy of  $425$  lm/W<sub>opt</sub> and a luminous efficiency of  $94$  lm/W, using near-UV LEDs with a 40% external quantum efficiency. Finally, we investigate FRET-converted light-emitting structures that use nonradiative energy transfer directly from epitaxial quantum wells to colloidal NCs. These proof-of-concept demonstrations show that FRET-based NC-LEDs hold promise for efficient solid-state lighting in green/yellow.

**Index Terms**—Förster resonance energy transfer (FRET), green/yellow, InGaN/GaN, LEDs, nanocrystals (NCs), nonradiative energy transfer.

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## I. INTRODUCTION

TODAY solid-state lighting is of great importance due to the financial benefits it offers (e.g., 50% saving in the global electricity consumption of lighting in principle) and its potential to combat environmental problems (e.g., climate change) [1]. In a recent study, the photometric conditions that achieve ultraefficient solid-state lighting have been investigated [2], [3]. According to this study, a luminous efficacy of optical radiation at  $408$  lm/W<sub>opt</sub> and a color rendering index (CRI) of 90 at a correlated color temperature (CCT) of 3000 K are simultaneously achievable. To obtain these photometric properties, LEDs that emit in blue, green, yellow, and red colors (at 463, 530, 573, and 614 nm with relative optical power levels of 1/8, 2/8, 2/8, and 3/8, respectively) are required [2], [3].

$\text{In}_x\text{Ga}_{1-x}\text{N}$  material system has the potential to cover the whole visible by changing the In composition ( $x$ ); however, it has technically been extremely challenging to obtain efficient green and yellow LEDs especially at these wavelengths (i.e., at 530 and 573 nm, respectively) [4]. Though blue-emitting InGaN/GaN LEDs may reach high external quantum efficiency [5], [6], those LEDs made of the same material system begin to suffer from significantly reduced internal quantum efficiency (and thus substantially reduced luminous efficiency) due to the increased In content in multiquantum wells (multi-QWs) while approaching green color [7]–[9]. In addition, the resulting optical emission is further reduced due to the increased internal fields in the QWs that separate electron and hole wavefunctions. Furthermore, undesired defects and compositional inhomogeneities also further decrease the efficiency of green LEDs [2].

For red LEDs, it is possible to achieve high quantum efficiencies around 650 nm by utilizing another material system, a quaternary alloy  $(\text{Al}_x\text{Ga}_{1-x})_{1-y}\text{In}_y\text{P}$ . By using this quaternary, it is possible to cover from 650 to 580 nm. However, the efficiencies significantly decrease approaching the yellow color. Therefore, there exists a significant gap in the green and yellow spectral regions to make efficient LEDs.

To address these problems, we propose and develop proof-of-concept green/yellow hybrid LEDs that use both radiative energy transfer and nonradiative energy transfer (Förster resonance energy transfer (FRET), or fluorescence resonance energy transfer) in their integrated colloidal semiconductor nanocrystals (NCs) for color conversion on near-UV LEDs (Fig. 1). Our experimental demonstrations indicate that the use of FRET

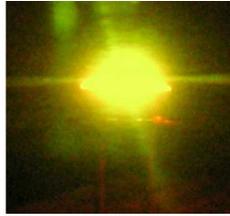


Fig. 1. Exemplary picture of a yellow hybrid LED while emitting at 573 nm.

in color conversion in green and yellow allows for significant enhancement in efficiency. Such FRET-converted LEDs hold great promise for efficient lighting, especially in green and yellow where the existing material systems and LED technologies suffer from low performance.

In conventional color conversion technique, phosphors are commonly used to generate white light. Phosphors provide wideband emission with high quantum efficiency across the visible. It is, however, difficult to tune their optical emission properties and obtain a spectrally narrow emission band (e.g., green or yellow emission alone). As an alternative, though, semiconductor NCs exhibit favorable properties for use in LED applications. They reveal tunable emission using quantum confinement effect [10], [11]. They show spectrally pure, narrow, and symmetric photoluminescence (PL) with large PL quantum yields and high photostability. They exhibit increasing absorption levels toward shorter wavelengths, relatively high photobleaching thresholds, and small spectral overlap between absorption and emission spectra [12]. Furthermore, various film deposition techniques including Langmuir–Blodgett, spin casting, and electrostatic layer-by-layer film deposition can be conveniently used to make good quality film. Therefore, these quasi-zero-dimensional semiconductor materials have been both theoretically and experimentally widely investigated [12]–[18]. Because of the tunable nature of their optical emission, they have also been exploited in various LED applications especially for hybrid white LEDs [19]–[24]. Unlike these previous studies of our group and the others, here we demonstrate hybrid LEDs that also employ FRET as a nonradiative means of energy transfer in color-converting colloidal NCs for increased efficiency in green/yellow spectral ranges for the first time.

In this paper, we present the design, growth, fabrication, integration, and characterization of NC-based LEDs that utilize direct and FRET-enhanced color conversion. Furthermore, we also investigate FRET-converted light-emitting structures. In our first NC-LED, we use green-emitting CdSe/ZnS core/shell NCs ( $\lambda_{PL} = 548$  nm) on near-UV InGaN/GaN LEDs ( $\lambda_{EL} = 379$  nm). In our second NC-LED that additionally employs FRET in NCs, we use a custom-design assembly of cyan and green-emitting CdSe/ZnS core/shell NCs ( $\lambda_{PL} = 490$  and 548 nm, respectively) to recycle the trapped excitons into NCs via FRET. In this proof-of-concept FRET-enhanced NC-LED, we achieve a quantum efficiency enhancement of 9% with respect to the first NC-LED. Finally, we examine FRET directly from epitaxial QWs to colloidal NCs in their close proximity.

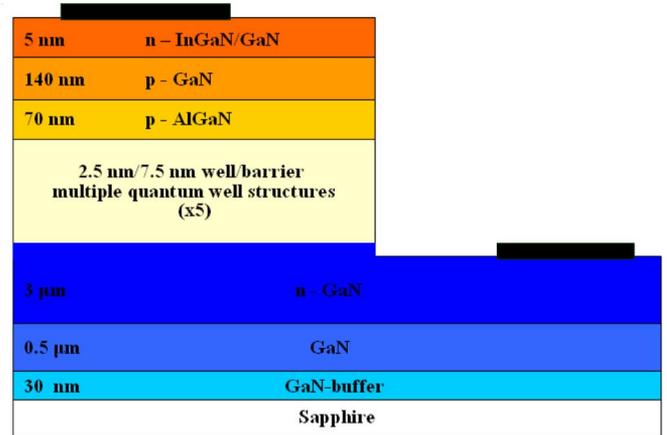


Fig. 2. Epitaxial design of our near-UV InGaIn/GaN LED.

Using FRET pumping in color conversion, a potential efficiency enhancement is predicted in such hybrid LEDs.

## II. NEAR-UV LED

For the growth of our near-UV LED, we employ a GaN dedicated metal–organic chemical vapor deposition (MOCVD) system and use a p–i–n epitaxial design for the near-UV LED, as shown in Fig. 2. At first, we begin with a 30-nm-thick GaN buffer layer and continue with a 0.5- $\mu\text{m}$ -thick GaN layer to increase the crystal quality of the device epitaxial layers. Subsequently, we grow a 3- $\mu\text{m}$ -thick Si-doped n-type contact layer. We continue with the epitaxial growth of 2.5 nm/7.5 nm InGaIn/GaN QW/barrier structure. We repeat the well/barrier structure five times for multi-QW region. Finally, we finish our growth with layers that consist of Mg-doped 70-nm-thick AlGaIn, 140-nm-thick Mg-doped GaN layer, and 2.5-nm-/2.5-m-thick Si-doped InGaIn/GaN layers. In the device fabrication, we use standard semiconductor processing including photolithography, thermal evaporator (metallization), reactive ion etch (RIE), and rapid thermal annealing. Finally, we make wire bonds to the device contacts in its package.

Our near-UV LED has a turn-ON voltage approximately at 3 V and an electroluminescence (EL) peak wavelength at 379 nm, as shown in the inset of Fig. 3. Our near-UV LED exhibits an external quantum efficiency of around 8% at low current levels. For example, the optical power is 5.8 mW at 19.05 mA. However, we observe a strong rollover in the efficiency with increasing input current, as our packaged LED does not include a heat sink and its full package is not designed for proper heat management. With a decreased level of device efficiency, the near-UV LED reaches an output optical power level of 8.5 mW at 43 mA, as shown in Fig. 3.

## III. COLOR-CONVERTED LED

For the first NC-LED, we use green-emitting CdSe/ZnS core/shell NCs as the color converters. Our green-emitting NCs have a concentration of 70 nmol/mL in toluene and their diameters are around 3.4 nm with a size dispersion of <5%. Their molecular weight is 100  $\mu\text{g}/\text{nmol}$  and they exhibit an in-solution

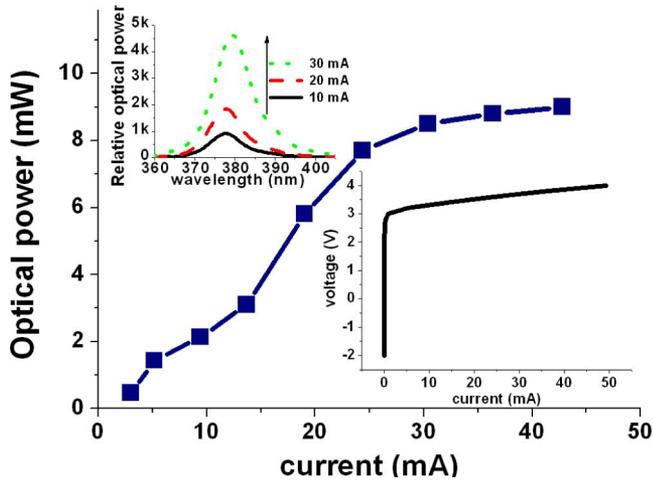


Fig. 3.  $L-I$  of our near-UV LED measured in an integrating sphere at room temperature (with no heat sink or active cooling) along with its EL spectra and  $I-V$  given in the insets.

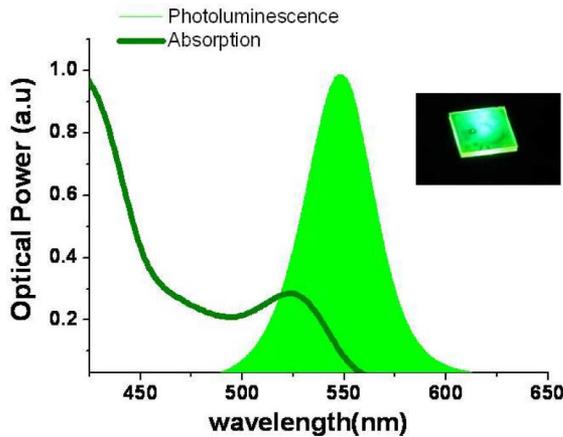


Fig. 4. Emission and absorption spectra of green-emitting CdSe/ZnS core/shell NCs along with their PL picture given in the inset.

quantum efficiency of  $>50\%$ . Their emission and absorption spectra are shown in Fig. 4. Toward shorter wavelengths, their optical absorption increases rapidly.

For hybrid light generation, we integrate closely packed NCs on a near-UV LED. The operating principle for such NC-LEDs relies on the optical excitation of the integrated NC films by the electrically driven near-UV LED. As a result of optical pumping, the NC PL generates the desired green/yellow light.

For our first NC-LED, we integrate a carefully determined amount of 9.1 nmol green-emitting CdSe/ZnS NCs on our LED. By using this exact amount of NCs, all of the LED EL at 379 nm is fully down-converted to NC emission by the NC luminophors. In Fig. 5, as the current injection level increases, the optical output power of our LED expectedly increases as well. When the current level reaches 19.05 mA, the optical power becomes 1.81 mW, which corresponds to a quantum efficiency of 46% for the integrated NC layer. The resultant luminescence of the hybrid LED is shown as a function of increasing current injection level in the inset of Fig. 5. Here, it is important to note that the green NC emission shifted 20 nm from 548 to 568 nm in

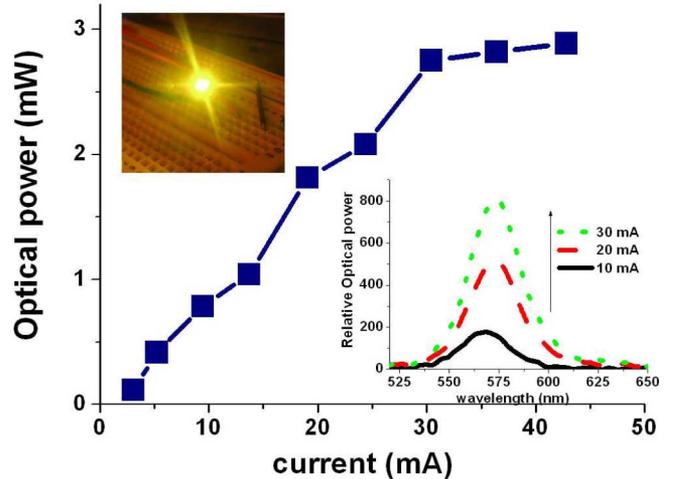


Fig. 5. Total optical power of green NC ( $\lambda_{PL} = 548$  nm) hybridized on near-UV LED ( $\lambda_{EL} = 379$  nm) at different levels of current injection at room temperature (with no heat sink or active cooling), along with the emission spectra and picture of the hybrid NC-LED when electrically driven.

TABLE I  
PHOTOMETRIC PROPERTIES OF OUR COLOR-CONVERTED NC-LED AT VARIOUS CURRENT INJECTION LEVELS

I (mA)	X	Y	CCT (K)
9.45	0.42	0.54	4008
19.50	0.45	0.50	3423
30.40	0.46	0.51	3260

the solid film. This is mainly due to the nonradiative energy transfer among NCs because of their size distribution ( $<5\%$ ). This size variance leads to slightly modified effective bandgaps. Consequently, the absorption of the relatively small-sized NCs and emission of the relatively large-sized NCs overlap, which enables energy transfer among them. However, this is not a fundamental problem. Since the starting NC peak emission can be conveniently set at any desired wavelength across the visible, this redshift can be easily compensated for using NCs with a shorter starting peak emission wavelength.

As the current injection level increases, the emission of packaged hybrid LED redshifts due to thermal effects that pushes our operating wavelength closer to 573 nm for yellow LED in ultra-efficient solid-state lighting. Thus,  $(x, y)$  tristimulus coordinates of the hybrid LED moves from (0.42, 0.54) to (0.46, 0.51) and its CCT shifts from 4008 to 3260 K, as summarized in Table I. We also measure the luminous efficacy of optical radiation and luminous efficiency of the source using (1) and (2), where  $v(\lambda)$  is the eye sensitivity function and  $P(\lambda)$  is the power spectral density (i.e., the light power emitted per unit wavelength) [6]. At 19.05 mA, our color-converted LED exhibits a luminous efficacy of 425 lm/W<sub>opt</sub> and a luminous efficiency of 17 lm/W

$$LE_{\text{Optical}} = \left( 683 \frac{\text{lm}}{\text{W}} \right) \frac{\int P_{\text{optical}}(\lambda)v(\lambda)d\lambda}{\int P_{\text{optical}}(\lambda)d\lambda} \quad (1)$$

$$LE_{\text{Electrical}} = \left( 683 \frac{\text{lm}}{\text{W}} \right) \frac{\int P_{\text{optical}}(\lambda)v(\lambda)d\lambda}{VI} \quad (2)$$

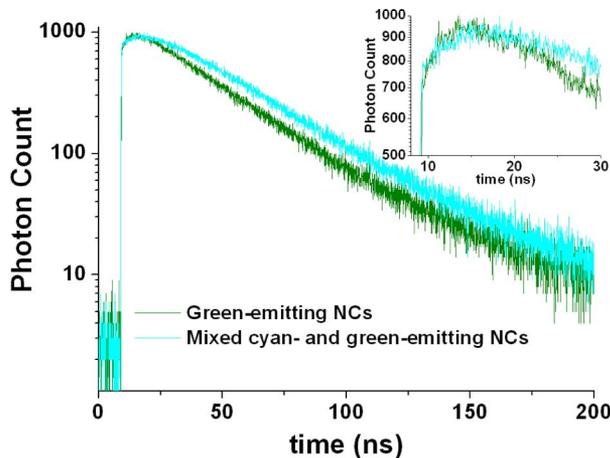


Fig. 6. Time-resolved spectroscopy measurements of the integrated green-emitting CdSe/ZnS core/shell NCs ( $\lambda_{PL} = 548$  nm) and mixed cyan- and green-emitting CdSe/ZnS core/shell NCs ( $\lambda_{PL} = 490$  and 548 nm, respectively) on near-UV LED ( $\lambda_{EL} = 379$  nm).

Fig. 6 shows the time-resolved spectroscopy measurements of the integrated NC layers at 556 nm. In Fig. 6, the energy rise (i.e., exponential increase) observed at 9 ns comes from the energy transfer. For these time-resolved spectroscopy measurements, we use a FluoTime 200 spectrometer (PicoQuant) with a time-correlated single photon counting (TCSPC) system of PicoHarp 300 that has a calibrated minimum possible time resolution of 4 ps. To pump NC solids, we use a laser head at an emission wavelength of 375 nm with light pulses as short as 70 ps and a photon multiplier tube (PMT) as the detector. Using an output monochromator, we characterize all of the prepared NC solids at 556 nm, corresponding to the peak emission wavelengths of the NCs. For the data analysis, we use the software FluoFit to take into account the instrumental response function (IRF).

To further analyze the PL decay curves of our NCs, we make fits to experimentally measured decay curves in Fig. 6 using a multiexponential model fit and convoluting with the laser diode response function at 375 nm. The resultant fitting parameters are summarized in Table II. Here, for our decays, we have two distinct lifetime components. The first lifetime component (associated with  $\tau_1$ ) is the general lifetime of the NCs, which is expectedly on the order of tens of nanoseconds [12], [13]. The second lifetime component (associated with  $\tau_2$ ) is an energy increase, which we conclusively differentiate with its negative amplitude. Here, it is worth noting that the distinct negative amplitude stems only possibly from the nonradiative energy transfer among NCs due to their size distribution, as discussed previously. For comparison, the cases of mixed green- and cyan-emitting NCs that exhibit intentionally strong energy transfer in Fig. 6 (and also in Table II) will be discussed in the next section.

#### IV. COLOR-CONVERTED FRET-ENHANCED LED

To enhance the efficiency of the hybrid LED, we integrate a carefully designed assembly of 3.4 nmol cyan-emitting NCs and 4.9 nmol green-emitting NCs as an energy gradient structure on a similar near-UV LED, which increases the overall quantum effi-

TABLE II  
MULTIEXPONENTIAL FITTING PARAMETERS FOR TRANSIENT PL DECAYS IN FIG. 6

	$A_1$	$\tau_1$ (ns)	$A_2$	$\tau_2$ (ns)
<b>Only green-emitting NCs</b>	219.69	32.78	-99.70	3.18
<b>Mixed cyan- and green-emitting NCs</b>	309.85	32.78	-166.77	12.10

$A_n$  and  $\tau_n$  are the associated amplitudes and decay time constants of multiexponential decay fits to the experimental PL decay curves;  $n = 1, 2$ .

ciency of the integrated NC layer. Here, the physical mechanism of the quantum efficiency enhancement stems from the ability to transfer excitation energy of the trapped excitons captured in defect states of donor NCs into acceptor NCs using nonradiative resonance energy transfer, as discussed in [25] and [26]. With the optical excitation by incoming near-UV light, the photo-generated electron and hole pairs relax rapidly (in hundreds of femtoseconds) to their respective conduction and valence bands, some of which subsequently make interband transitions for radiative recombination. However, some of these photogenerated electron-hole pairs get trapped in the defect states, typically leading to nonradiative recombination. Using nonradiative energy transfer, some portions of these trapped excitons are transferred to the green-emitting NCs that serve as acceptors for these defect states, a part of which, in turn, contribute to PL of these acceptors. Thus, the overall luminescence increases because of the enhanced emission of the acceptors by effectively recycling the trapped excitons (and thus otherwise lost excitation energy) in the defect states of donors.

In Fig. 6, the time-resolved spectroscopy measurements of mixed cyan- and green-emitting NC film are presented. We observe in Fig. 6 (and also in Table II) that the energy rise (i.e., the exponential increase) in the energy gradient structure of mixed cyan-green NCs is more dominant with respect to the case of only green NCs, because of the strong exciton migration to green-emitting NCs. By using the physical process of FRET, we obtain the resulting luminescence of the hybrid device, as shown in Fig. 7. Here, the quantum efficiency of the NC film is enhanced by 9%.

It is important to note that cyan-emitting NCs are fully quenched by transferring their energy to green-emitting NCs. As a result of the enhanced quantum efficiency of the mixed NC assembly, at 19.05 mA, the FRET-enhanced NC-LED reaches a luminous efficiency of 19 lm/W, while the luminous efficacy of optical radiation is maintained at 425 lm/W<sub>opt</sub>. This FRET-NC-LED has its chromaticity coordinates at  $(x, y) = (0.46, 0.51)$ . As the current injection level increases, the emission of FRET-NC-LED makes expectedly a redshift from 568 to 573 nm due to the additional heating. Its tristimulus coordinates and CCT correspondingly shift from (0.40, 0.49) to (0.45, 0.52) and from 4077 to 3441 K, respectively, as given in Table III. In our FRET-NC-LED, NCs reach a quantum efficiency level as high as 55%. By using a state-of-the-art near-UV LED with an external quantum efficiency of 40%, this implies that it is possible to obtain color-converted FRET-enhanced LEDs with an expected external quantum efficiency of 22% and a predicted luminous efficiency of 94 lm/W.

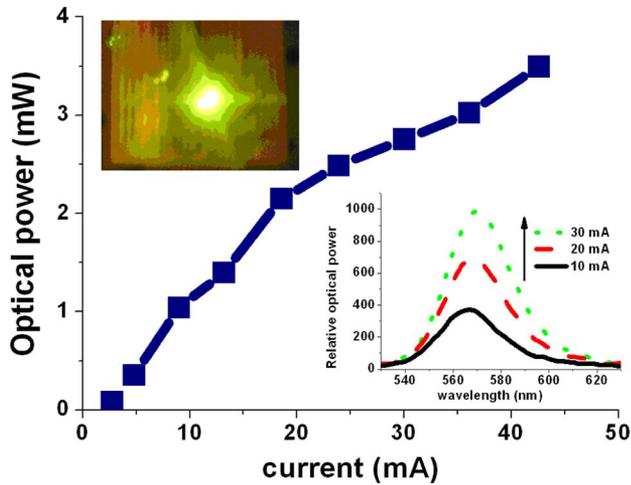


Fig. 7. Total optical power of mixed cyan- and green-emitting NCs ( $\lambda_{PL} = 490$  and  $548$  nm, respectively) hybridized on near-UV LED ( $\lambda_{EL} = 379$  nm) at different levels of current injection at room temperature (with no heat sink or active cooling), along with the emission spectra and picture of the resulting FRET-NC-LED when electrically driven.

TABLE III  
PHOTOMETRIC PROPERTIES OF OUR COLOR-CONVERTED FRET-ENHANCED NC-LED AT VARIOUS CURRENT INJECTION LEVELS

I (mA)	x	Y	CCT (K)
9.45	0.40	0.49	4077
19.50	0.44	0.51	3642
30.40	0.45	0.52	3441

V. FRET-CONVERTED LIGHT-EMITTING STRUCTURES

In the first NC-LED, color conversion is achieved through only optical pumping (i.e., radiative energy transfer) that involves two recombination processes. The first radiative recombination process occurs in the electrically driven QWs of the LED, and subsequently, their emitted photons excite the integrated NC luminophors that further luminescence via a second radiative recombination process. On the other hand, rather than using two-step recombination processes, color conversion using only one recombination step is also possible via nonradiative energy transfer directly from epitaxial QWs to colloidal NCs [which are chemically synthesized quantum dots (QDs)] in their close proximity [13]. The advantage of such energy-transferring hybrid structure is that recombination takes place directly in the NCs (i.e., after QWs transfer their excitation energy to these QDs for the recombination process to directly take place in the dots) for color conversion. Therefore, using nonradiative energy transfer for color conversion facilitates a direct excitation of NCs with a more efficient, single-step process, as opposed to merely optical pumping that requires two steps of radiative recombination.

To nonradiatively pump the NCs by using FRET, we use a special epiwafer design, as shown in Fig. 8. In its epitaxial growth, we start with a 30-nm-thick GaN buffer layer on a double-side-polished c-plane sapphire and a 2- $\mu$ m-thick GaN buffer layer to increase the crystal quality of the device epitaxial layers. Subsequently, we grow a 1.2- $\mu$ m-thick Mg-doped p-type

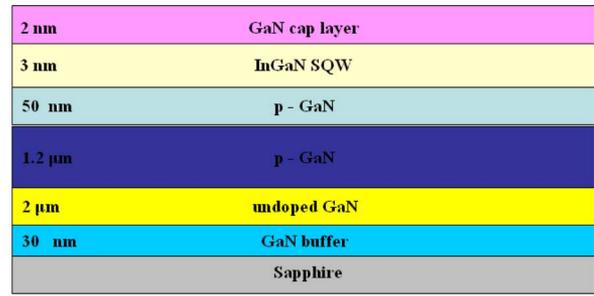


Fig. 8. Epitaxial structure of our wafer for energy transfer from single QW (SQW) to colloidal NCs for yellow/green emission.

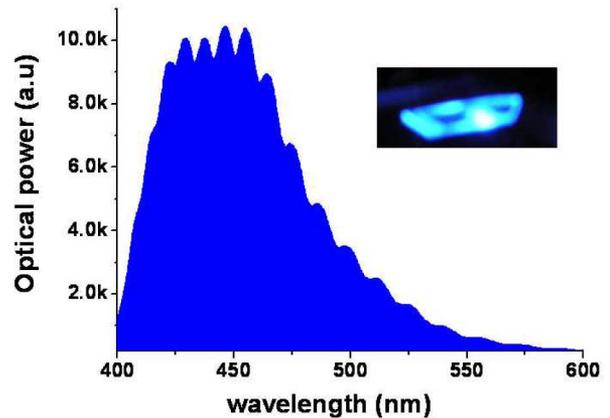


Fig. 9. PL of our epiwafer for energy transfer from SQW to colloidal NCs.

GaN layer and a 50-nm-thick highly doped p-type GaN layer. We then continue the epi-growth with a 3-nm-thick InGaN well. Finally, we finish our growth with a 2-nm cap layer that enables the dipole–dipole interaction of the QW and NCs. For steady-state measurements, we use a Jobin Yvon Triax 500 CCD PL system with an He/Cd laser at 325 nm, and the PL of the resulting wafer is shown in Fig. 9. The resultant wavy PL behavior stems for phase separation of the QW. Here, since the optical absorption of the NCs increases toward shorter wavelengths (as shown in the inset of Fig. 4), the overlap of the QW emission and the NC absorption makes strong dipole–dipole interaction between these NCs and QW possible.

In this color-converting hybrid structure, the energy transfer is possible as a result of the dipole–dipole interaction between the QW and NCs, rather than the delocalized behavior of excitons or free carriers, because ZnS barriers of our NCs prevent tunneling and provide full electronic isolation. We prepare our hybrid structure by integrating NCs on top of the epiwafer and its negative control group by integrating the same density and amount of NCs on quartz substrate. To explore the energy transfer dynamics, we investigate the time-resolved fluorescence of our hybrid system at 556 nm to observe the energy feeding into NCs, instead of studying the quenching of the QWs. This is because the quenching can also occur because of environmental trap states, but the energy feeding in NC solids can occur merely as a result of the energy transfer to them.

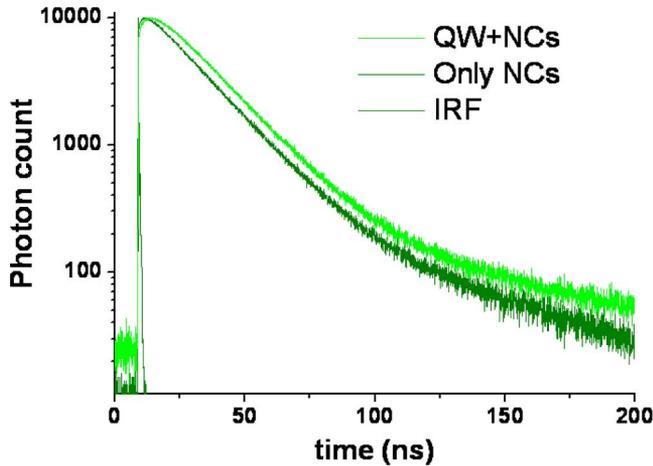


Fig. 10. Time resolved spectroscopy measurements of our hybrid FRET-converted system that consists of both QW and NCs (QW+NCs), compared with that of the negative control group that contains only NCs (only NCs), along with the instrument response function (IRF) of the laser diode emitting at 375 nm, using a TCSPC system of PicoHarp 300 with a time resolution of 64 ps.

TABLE IV  
MULTIEXPONENTIAL FITTING PARAMETERS FOR TRANSIENT PL DECAYS  
IN FIG. 10

	$A_1$	$\tau_1$ (ns)	$A_2$	$\tau_2$ (ns)	$A_3$	$\tau_3$ (ns)
Only green-emitting NCs	14339	18.34	351	70.44	-4244	4.36
QWs+ green-emitting NCs	21651	18.34	695	61.43	-13796	7.29

$A_n$  and  $\tau_n$  are the associated amplitudes and decay time constants of multiexponential decay fits to the experimental PL decay curves;  $n = 1, 2, 3$ .

Fig. 10 shows the time-resolved spectroscopy of our hybrid system and that of only the NC solids as the control group. Table IV also summarizes the corresponding multiexponential fitting parameters for these transient PL decays. Here, the first lifetime component (associated with  $\tau_1$ ) is again the general lifetime of NC emitters on the order of tens of nanoseconds. The second lifetime component (associated with  $\tau_2$ ) is relatively weak in amplitude and is attributed to the environmental effects. The third lifetime component (associated with  $\tau_3$ ) distinctively exhibits negative amplitude, indicating nonradiative energy transfer. In the hybrid FRET-converted structure, this third lifetime component is increased, leading to emission decaying further slower. In agreement with this observation, the transient PL in Fig. 10 also shows a smoother turnaround from the rising edge to the decaying tail (from 8 to 10 ns) owing to the energy transfer feeding component of the resulting PL.

In Fig. 11, the steady-state PL spectra of our hybrid FRET-converted system that contains NCs integrated on QW and that of its negative control group that includes only the NCs are depicted. Here, the emission of the NCs in the hybrid system is increased with respect to the case of only NCs as a result of the energy transfer from QW to NCs. Furthermore, in the hybrid system, the emission of NCs is spectrally disturbed, resulting in an asymmetric spectral emission profile, which is another evidence of the nonradiative energy transfer from the QW to NCs.

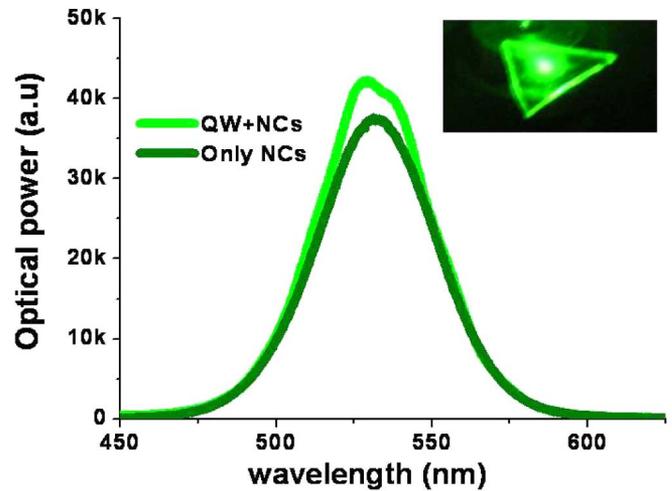


Fig. 11. Steady-state emission spectrum of our hybrid FRET-converted system (QW + NCs) improved by using nonradiative energy transfer pumping of green-emitting CdSe/ZnS core/shell NCs ( $\lambda_{PL} = 548$  nm) by blue-emitting InGaN/GaN QW ( $\lambda_{PL} = 450$  nm) with respect to the case of only NCs.

Nonradiative energy transfer pumping from QWs to NCs for color conversion allows for direct excitation of the NC luminophors, leading to a subsequent single-step radiative recombination process, compared to only radiative energy transfer pumping that requires two subsequent steps of radiative recombination. Using nonradiative energy transfer in addition to radiative transfer, it is possible to make color-converted FRET-enhanced LEDs with an external quantum efficiency of 22%, starting with a state-of-the-art near-UV LED that has an external quantum efficiency of 40%. In the case of using only nonradiative energy transfer pumping, though, it is predicted to possibly improve the overall efficiency even further. These preliminary results on FRET conversion implies that such NC-based FRET-converted solid-state lighting holds great promise for future lighting applications.

## VI. CONCLUSION

In this study, we presented the design, growth, fabrication, hybridization, and characterization of hybrid LEDs that integrate colloidal semiconductor NCs as color converters on near-UV InGaN/GaN LEDs. In our proof-of-concept color-converted NC-LED by hybridizing green-emitting CdSe/ZnS core/shell NCs, we achieved the luminous efficacy of optical radiation to be 425 lm/W<sub>opt</sub> and the luminous efficiency to be 17 lm/W at a current injection level of 19.05 mA.

In our proof-of-concept color-converted FRET-enhanced NC-LED by hybridizing a custom-design assembly of cyan- and green-emitting CdSe/ZnS core/shell NCs to facilitate nonradiative energy transfer among them, we demonstrated a quantum efficiency enhancement of 9% with respect to the color-converted NC-LED by recycling the trapped excitons into NCs via FRET. At a current injection level of 19.05 mA while keeping the luminous efficacy of optical radiation at 425 lm/W<sub>opt</sub>, we further improved our hybrid NC-LED to reach a luminous efficiency of 19 lm/W. In our FRET-NC-LED, NCs reach a quantum

efficiency level of 55%. Starting with a 40% efficiency near-UV LED, a luminous efficiency of 94 lm/W and an external quantum efficiency of 22% are predicted in yellow/green spectral ranges.

Finally, investigating FRET-converted light-emitting structures for direct excitation from epitaxial QWs to colloidal NCs, a further potential efficiency enhancement is predicted for such hybrid LEDs. Currently, experiments on the quantification of FRET conversion enhancement are underway. These proof-of-principle demonstrations show that these hybrid NC-LEDs are potentially strong candidates for future lighting applications in green/yellow.

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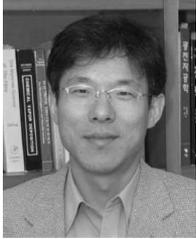
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